

See-saw tau lepton mass and calculable neutrino masses in a 3-3-1 model

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In this work we show that in a version of the 3-3-1 model proposed by Duong and Ma, in which the introduction of a scalar sextet is avoided by adding a singlet heavy charged lepton, the τ lepton gains mass through a see-saw-like mechanism. We also show how to generate neutrino masses at the 1-loop level, and give the respective Maki-Nakagawa-Sakata mixing matrices for a set of the parameters. We also consider the effect of adding a singlet right-handed neutrino.

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Some years ago it was proposed a model with $SU(3)_C \otimes SU(3)_L \otimes U(1)_N$ gauge symmetry, 3-3-1 for short, in which the three lepton families transform in the same way under the gauge symmetry i.e., $\Psi_{aL} = (\nu_a, l_a, l_a^c)_L \sim (\mathbf{1}, \mathbf{3}, 0)$ with $a = e, \mu, \tau$ [1,2]. Therefore, in this model the lepton mass term transforms as $(\mathbf{1}, \mathbf{3}, 0) \otimes (\mathbf{1}, \mathbf{3}, 0) = (\mathbf{1}, \mathbf{3}^*, 0)_A \oplus (\mathbf{1}, \mathbf{6}, 0)_S$, hence, in order to give mass to the charged leptons it is possible to introduce a triplet $\eta = (\eta^0, \eta_1^-, \eta_2^+)^T \sim (\mathbf{1}, \mathbf{3}, 0)$ and a symmetric sextet $S \sim (\mathbf{1}, \mathbf{6}, 0)$ [2]. With the η -triplet only, one of the charged leptons remains massless and the other two are mass degenerated. Hence, at least a sextet S has to be introduced in order to give arbitrary masses to all charged leptons. Although this implies that the model has a rather complicated Higgs scalar sector, it is interesting to stress that all the extensions of the electroweak standard model with extra Higgs scalar multiplets, like multi-Higgs doublets [3], singlets [4,5], and triplets [6], are embedded in this 3-3-1 model. In fact, under the subgroup $SU(2)_L \otimes U(1)_Y \subset SU(3)_L \otimes U(1)_N$ the model has three scalar doublets, four singlets (one neutral, one singly charged and two doubly charged), and a complex triplet. This sort of models also gives some insight on the families replication and on the observed value of the weak mixing angle $\sin^2 \theta_W < 1/4$. Moreover, recently the 3-3-1 models are interesting possibilities for explaining the new value of the positive muon anomalous magnetic moment, reported by the Muon ($g - 2$) Collaboration [7]. The new ($g - 2$) value is consistent not only with supersymmetry [8] but also with several versions of the 3-3-1 model as it was shown in Ref. [9]. Finally, in a version of

the model it is possible to generate the top and bottom quark masses at the tree level while the other quarks and charged leptons gain mass at the 1-loop level [10].

The see-saw mechanism was proposed in order to understand the smallness of the neutrino masses [11]. On the other hand, later, it was introduced a generalization of this mechanism, which is also valid for the charged lepton masses, in the context of the minimal left-right symmetric model [12,13]. Here we will show, in a 3-3-1 model, that a similar mechanism can be implemented for the case of the τ lepton mass and that the three neutrinos gain mass radiatively. In the case of three neutrinos the masses and mixing are almost completely determined by the charged lepton parameters. In fact they depend only upon a coupling constant (λ) of a quartic term of the scalar potential. It implies that we are also calculating the ratio of the neutrino masses (see Eq. (12) below).

Some years ago, Duong and Ma [14] noted that a way to give mass to the charged leptons that does not involve a scalar sextet can be implemented: Besides the usual scalar triplets: η , mentioned before, $\rho = (\rho^+, \rho^0, \rho^{++})^T \sim (\mathbf{1}, \mathbf{3}, 1)$ and $\chi = (\chi^-, \chi^{--}, \chi^0)^T \sim (\mathbf{1}, \mathbf{3}, -1)$, they introduced a charged lepton transforming like singlet: $E'_L \sim (\mathbf{1}, \mathbf{1}, -1)$ and $E'_R \sim (\mathbf{1}, \mathbf{1}, -1)$. In this case the Yukawa interactions are given by

$$\begin{aligned} -\mathcal{L}_Y = & \frac{1}{2} \sum_{\substack{a,b=e,\mu,\tau \\ i,j,k=1,2,3}} \epsilon^{ijk} \overline{(\Psi_{aLi})^c} F_{ab}(\Psi)_{bLj} \eta_k \\ & + \sum_{a=e,\mu,\tau} [f_a \overline{(\Psi)}_{aL} E'_R \rho + f'_a \overline{E'_L} \chi^T (\Psi^c)_{aR}] \end{aligned}$$

$$+ M \bar{E}'_L E'_R + H.c., \quad (1)$$

where F_{ab} denotes an antisymmetric matrix. Then the mass matrix for the charged leptons, in the basis of the symmetry eigenstates that we denote $l'_{L,R} = (e', \mu', \tau', E')_{L,R}$, reads $\bar{l}'_L \mathcal{M}^l l'_R + H.c.$ where

$$\mathcal{M}^l = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & -f_{e\mu}v_\eta & -f_{e\tau}v_\eta & f_e v_\rho \\ f_{e\mu}v_\eta & 0 & -f_{\mu\tau}v_\eta & f_\mu v_\rho \\ f_{e\tau}v_\eta & f_{\mu\tau}v_\eta & 0 & f_\tau v_\rho \\ f_e v_\chi & f'_\mu v_\chi & f'_\tau v_\chi & \sqrt{2}M \end{pmatrix}. \quad (2)$$

Here for simplicity we have assumed real vacuum expectation values: $\langle \eta^0 \rangle \equiv v_\eta/\sqrt{2}$, $\langle \rho^0 \rangle \equiv v_\rho/\sqrt{2}$ and $\langle \chi^0 \rangle \equiv v_\chi/\sqrt{2}$.

Firstly, by assuming that the only vanishing dimensionless parameters in Eq. (2) are $f_{e\mu}, f_{e\tau}, f_e$ and f'_e , we obtain $m_e = 0$ and approximately

$$\begin{aligned} m_\mu &\approx -\frac{f_\mu f_{\mu\tau}}{\sqrt{2}f_\tau} v_\eta, & m_\tau &\approx -\frac{f_\tau f'_\tau v_\rho v_\chi}{2M} - m_\mu, \\ m_E &\approx M - m_\tau + m_\mu. \end{aligned} \quad (3)$$

To get a positive mass for the μ and τ leptons we can redefine the respective field by a γ_5 factor i.e., $\mu(\tau) \rightarrow \gamma_5 \mu(\tau)$ (or see below). We see from Eq. (3) that the τ lepton mass is of the see-saw type [12,13], however, as showed below, the M mass is not related to a grand unification scale.

The mass matrix in Eq. (2) is diagonalized by a biunitary transformation $U_L^{l\dagger} \mathcal{M}^l U_R^l = \hat{\mathcal{M}}^l$, relating the symmetry eigenstates (primed fields) with the mass eigenstates (unprimed fields), $l'_L = U_L^l l_L$ and $l'_R = U_R^l l_R$, where the mass eigenstates are denoted by $l_{L,R} = (e, \mu, \tau, E)_{L,R}$.

Notice that even for massless neutrinos, which may result in diagonal interactions with the W boson, there are flavor changing neutral and charged currents in the interactions in Eq. (1). It means that although the model has many parameters they are not fixed only by the mass of the charged leptons but by new interactions too. In this situation, we can introduce the unitary matrices $U_{L,R}^l$ in Eq. (1) and leaves that phenomenology determine both, the masses of the Higgs scalars and the matrix elements of those matrices or, we can determine first those matrices by given the appropriate mass to the charged leptons and use phenomenology to determine only the mass of the Higgs scalars. Here we will make the first part of the second alternative, i.e., we will determine the mixing matrices $U_{L,R}^l$ and left for further studies the phenomenology of the model.

Since we will be concerned below with the neutrino mass phenomenology, we find appropriate to obtain a set of parameters which are consistent with the actual values of the charged lepton masses. In this vain we use $U_L^{l\dagger} \mathcal{M}^l (\mathcal{M}^l)^\dagger U_L^l = (\hat{\mathcal{M}}^l)^2$ and $U_R^{l\dagger} (\mathcal{M}^l)^\dagger \mathcal{M}^l U_R^l = (\hat{\mathcal{M}}^l)^2$ in order to find $U_{L,R}^l$. The numerical analysis of the

eigenvalues of $\mathcal{M}^l (\mathcal{M}^l)^\dagger$ (or $(\mathcal{M}^l)^\dagger \mathcal{M}^l$), gives the following masses for the charged leptons (in GeV):

$$\begin{aligned} m_e &= 0.000510985, & m_\mu &= 0.105658, \\ m_\tau &= 1.77703, & m_E &= 3007.79, \end{aligned} \quad (4)$$

which are in good agreement with those of Ref. [15], if the following parameters are chosen:

$$\begin{aligned} f_{e\mu} &= f_{e\tau} = 0.0006785, & f_{\mu\tau} &= 0.021, \\ f_e &= 0.0004263, & f'_e &= 0.001, & f_\mu &= 0.0481815, & f'_\mu &= 0.05, \\ f_\tau &= 0.1259848, & f'_\tau &= 0.3; \end{aligned} \quad (5)$$

and also $v_\chi = 1000$ GeV, $M = 3000$ GeV, $v_\eta = 20$ GeV and $v_\rho = 245.186 (= \sqrt{246^2 - 20^2})$ GeV. The left- and right-handed mixing matrices, up to five decimal places, are given by

$$U_L^l \sim \begin{pmatrix} -0.99517 & -0.09776 & 0.00836 & 0.00002 \\ 0.08889 & -0.86222 & 0.49866 & 0.00276 \\ -0.04154 & 0.49701 & 0.86672 & 0.00724 \\ 0.00008 & -0.00122 & -0.00766 & 0.99997 \end{pmatrix}, \quad (6)$$

$$\tilde{U}_R^l \sim \begin{pmatrix} 0.99914 & -0.04130 & 0.00433 & 0.00024 \\ 0.04133 & 0.99905 & -0.00754 & 0.01175 \\ 0.00395 & -0.00852 & -0.99747 & 0.07053 \\ -0.00100 & -0.01116 & 0.07061 & 0.99744 \end{pmatrix}. \quad (7)$$

We have verified that in fact, with the matrices given in Eqs. (6) and (7) with $U_L^{l\dagger} \mathcal{M}^l \tilde{U}_R^l = \hat{\mathcal{M}}^l$ where $\tilde{U}_R^l = U_R^l \phi$ with $\phi = \text{diag}(1, -1, 1, -1)$ and $\hat{\mathcal{M}}^l = \text{diag}(m_e, m_\mu, m_\tau, m_E)$. This analysis illustrate how the Duong and Ma singlet lepton works [14] to give the correct mass for the charged leptons.

Until this point neutrinos remain massless. However, there are several ways to give mass to the neutrinos in the context of the 3-3-1 models. These include the mass generation at the tree level [16] or by radiative corrections [17]. Here neutrinos will continue to be massless unless we allow the breakdown of the $U(1)_{B+L}$ global symmetry in the scalar potential, where L is the total lepton number and B is the baryon number [18,19]. This is because if $B+L$ is conserved, the mixture in the singly charged scalar sector occurs only between η_1^-, ρ^- , and between η_2^-, χ^- , and diagrams like that in Fig. 1 do not occur, see Eq.(8). Recall that the $B+L$ assignments are 2 for η_2^- and χ^- and 0 for η_1^- and ρ^- [19], for leptons $B+L$ coincide with L or with the family lepton number. However if $B+L$ is not conserved a quartic term $\lambda(\eta^\dagger \chi)(\eta^\dagger \rho) + H.c.$ is allowed and all singly charged scalar fields mix one to another. We will consider the case in which this quartic coupling among the scalar Higgs does exist and diagrams like that in Fig. 1 do occur. The charged current is written in the mass-eigenstate basis as

$\bar{l}_L \gamma^\mu V_{MNS} \nu_L W_\mu^-$ with the Maki-Nakagawa-Sakata matrix [20] defined as $V_{MNS} = (O_L^l)^\dagger U_L^\nu$ where U_L^ν is the matrix which relates the symmetry and the left-handed neutrino mass eigenstates. Notice that in this case V_{MNS} is a 4×3 matrix since O_L^l is the submatrix 3×4 of U_L^l , i.e., omitting the forth row in Eq. (6).

Let us now turn on the lepton interactions with the triplet η : From Eq. (1) we have

$$-\overline{l}_{\alpha R} (O_R^{lT})_{\alpha a} F_{ab} \nu'_L \eta_1^- - \overline{\nu'^c_R} F_{ab} (O_L^l)_{ba} l_{\alpha L} \eta_2^+. \quad (8)$$

Notice that also in Eq. (8) only the 3×4 submatrix of $U_{L,R}^l$ appear, i.e., omitting the forth row in Eq. (6) and (7), since $a, b = e, \mu, \tau$ but $\alpha = e, \mu, \tau, E$; we have already denoted those submatrices by O_L^l, O_R^l .

Hence, we have a finite 3×3 non-symmetric neutrino mass matrix at the 1-loop level of the form

$$\begin{aligned} \mathcal{M}_{bd}^\nu &= \lambda v_\rho v_\chi \sum_{\alpha, a, f} (O_R^{lT})_{\alpha a} F_{ab} m_\alpha F_{df} (O_L^l)_{fa} I(M_1^2, M_2^2, m_\alpha^2) \\ &\approx \frac{\lambda v_\rho v_\chi}{16\pi^2 M_2^2} \sum_\alpha m_\alpha (O_R^{lT} F)_{\alpha b} (F O_L^l)_{d\alpha} I(r, s_\alpha), \end{aligned} \quad (9)$$

and

$$I(r, s_\alpha) = \int_0^1 dx \int_0^{1-x} dy \frac{1}{(s_\alpha^2 - r^2)x + (s_\alpha^2 - 1)y - s_\alpha^2}, \quad (10)$$

where $M_{1,2}$ denote typical masses for the singly charged scalars and we have defined $r = M_1/M_2$ and $s_\alpha = m_\alpha/M_2$. We have neglected a similar contribution proportional to v_η^2 , which arises from Fig. 1 if we make $\eta_1^- \rightarrow \rho^-$, $\eta_2^- \rightarrow \chi^-$ and $v_\rho v_\chi \rightarrow v_\eta^2$, since we are assuming that $v_\eta^2 \ll v_\rho v_\chi$.

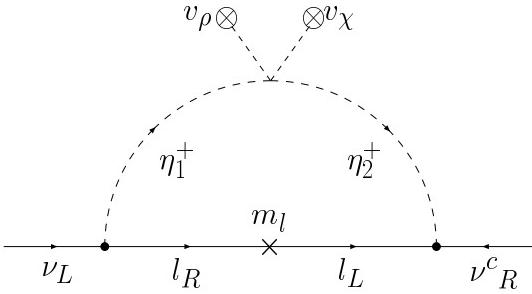


FIG. 1. Diagram showing calculable neutrino masses.

From Eq.(9) we can built the symmetric mass matrix $M^\nu = (\mathcal{M}^\nu + \mathcal{M}^{\nu T})/2$. With $M_1 = 100$ GeV, $M_2 = 1$ TeV, and the parameters already obtained for the charged leptons in Eq. (5), we have from Eqs. (9) and (10) numerical values for M^ν , in GeV and up to a factor 10^{-6} ,

$$M^\nu \sim \lambda \begin{pmatrix} -0.00775 & -0.18972 & 0.05030 \\ & -4.60084 & 1.2635 \\ & & -0.29346 \end{pmatrix}. \quad (11)$$

This matrix has the following mass eigenvalues:

$$\lambda (\approx 0, 0.04931, 4.95892) 10^{-6} \text{ GeV}, \quad (12)$$

where the first eigenvalue is assumed to be zero up to the decimal places we are considering.

This time we diagonalize the mass matrix in Eq. (11) by making $\Phi U^{\nu T} M^\nu U^\nu \Phi = \hat{M}^\nu = \text{diag}(m_1, m_2, m_3)$ with $m_3 > m_2 > m_1$ where $\Phi = \text{diag}(1, 1, i)$ and

$$U^\nu \sim \begin{pmatrix} -0.99896 & -0.02268 & 0.03961 \\ 0.03228 & 0.26260 & 0.96436 \\ -0.03228 & 0.96464 & -0.26160 \end{pmatrix}. \quad (13)$$

Notice that the U^ν matrix does not depend on the global λ parameter. In this case we have the mixing matrix $V_{MNS} \equiv (O_L^l)^T U^\nu \Phi$:

$$V_{MNS} \sim \begin{pmatrix} 0.9983 & 0.0058 & 0.0572i \\ 0.0538 & 0.2552 & -0.9654i \\ -0.0202 & 0.9668 & 0.2545i \\ -0.0002 & 0.0077 & 0.0008i \end{pmatrix}, \quad (14)$$

where we have used the basis $l = (e, \mu, \tau, E)^T$ and $\nu = (\nu_1, \nu_2, \nu_3)^T$. Although the first row and column of this mixing matrix is compatible with the SMA mixing matrix, we cannot compare directly these numbers with the usual ones [21–23] because the matrix in Eq. (14) is not unitary, it means that, in the case of three neutrinos (three masses), there are more parameters than in the case of an unitary mixing matrix (three Euler angles).

Now we can try to fix the trilinear λ parameter to be consistent with the neutrino mass phenomenology: From atmospheric neutrino data [24] we know that at least one mass eigenstate does exist such that its mass is $m > \sqrt{|\Delta m_{\text{atm}}^2|} \sim (4 - 6) \cdot 10^{-2}$ eV [25]. On the other hand, the solar neutrino data requires $\Delta m_\odot^2 \sim 10^{-5}, 10^{-7}, 10^{-10}$ eV² depending of the type of solution LM and SM, LOW or VO solution, respectively [26]. Hence by coherence we will choose λ in order to be consistent with the SMA scenario. From Eq. (12) we see that with $\lambda \sim 6.4 \times 10^{-5}$ we obtain $m_1 = 0$ eV, $m_2 = 3.2 \cdot 10^{-3}$ eV and $m_3 = 3.2 \cdot 10^{-1}$ eV.

Moreover, we would like to stress that the values of the parameters given in Eq. (5) are not expected to be unique and it means that it could be possible to obtain other values in Eq. (11) and (12) and of course, of the MNS-matrix in Eq. (14), which could fit better the experimental data [21]. We have verified also that the neutrino masses and the mixing matrix do not strongly depend on the masses of the charged scalars M_1, M_2 . In this way, the neutrino parameters are a prediction of the model: Once the f 's constants in Eq. (2) are made consistent with the known charged lepton masses, the ratios of the neutrino masses and the mixing matrix U^ν are also automatically fixed, since λ is a common factor.

Next, we would like to point out that in the 3-3-1 model the introduction of fermion singlets is more natural than in other models, including the standard electroweak model, because in the present context the charged singlet

is introduced in order to replace the scalar sextet. Since the number of fermion families transforming non-trivially under the 3-3-1 gauge symmetry must be a multiple of 3, the more economic way to introduce new fermions is by adding singlets under this symmetry. Moreover, in the present model which has already a charged lepton singlet, if we want to have again a symmetry in the representation content between the charged leptons and neutrinos we can add a right-handed neutrino singlet N'_R . In this case we can obtain a more wide neutrino mass spectrum if we consider N'_R as a light neutrino. In this situation we have the interactions

$$\sum_{a=e,\mu,\tau} \hat{h}'_a \overline{\psi_a L} N'_R \eta + \frac{1}{2} \overline{(N'_R)^c} M_R N'_R + H.c. \quad (15)$$

The neutrino mass matrix in the basis $N' = (\nu'_{eL}, \nu'_{\mu L}, \nu'_{\tau L}, (N'^c)_L)^T$ and up to a 10^{-6} factor, in GeV, is given by

$$M^\nu \sim \lambda \begin{pmatrix} -0.00775 & -0.18972 & 0.05030 & \frac{\hat{h}_e v_\eta}{2\sqrt{2}\lambda} \\ & -4.60838 & 1.2635 & \frac{\hat{h}_\mu v_\eta}{2\sqrt{2}\lambda} \\ & & -0.29346 & \frac{\hat{h}_\tau v_\eta}{2\sqrt{2}\lambda} \\ \frac{\hat{h}_e v_\eta}{2\sqrt{2}\lambda} & \frac{\hat{h}_\mu v_\eta}{2\sqrt{2}\lambda} & \frac{\hat{h}_\tau v_\eta}{2\sqrt{2}\lambda} & \frac{M_R}{2\lambda} \end{pmatrix}, \quad (16)$$

where the numerical 3×3 submatrix denotes the mass entries generated by radiative corrections already obtained in Eq. (11); we have defined $\hat{h}_a = 10^6 \hat{h}'_a$. The mass matrix in Eq. (16) has the following mass eigenvalues:

$$(\approx 0, 0.005, 0.052, 1.052) \text{ eV}, \quad (17)$$

if we chose the following values for the dimensionless parameters

$$(\lambda, \hat{h}_e, \hat{h}_\mu, \hat{h}_\tau) = (10^{-6}, 2 \cdot 10^{-3}, 10^{-7}, 10^{-7}), \quad (18)$$

and $v_\eta = 20$ GeV, $M_R = 2$ eV. We can identify, for instance, $\Delta m_{43}^2 = 1.1 \text{ eV}^2$, $\Delta m_{32}^2 = 2.7 \times 10^{-3} \text{ eV}^2$ and $\Delta m_{21}^2 = 2.1 \times 10^{-5} \text{ eV}^2$ with the respective mass square difference consistent with LSND [27], atmospheric [24], and solar [26] neutrino data. In this case the mixing matrix in the neutrino sector is given by

$$U^\nu \sim \begin{pmatrix} 0.99912 & -0.04199 & -0.00191 & -0.00067 \\ -0.04193 & -0.99886 & 0.02275 & -0.00091 \\ -0.00293 & -0.02191 & -0.97602 & -0.21654 \\ 0 & 0.00582 & 0.21646 & -0.97627 \end{pmatrix}. \quad (19)$$

The relation between phenomenological neutrinos N'_a and the massive ones $N = (\nu_1, \nu_2, \nu_3, \nu_4)^T$ is given by $N' = U^\nu \Phi N$ where $\Phi = \text{diag}(1, i, i, 1)$. The $V_{MNS} = U_L^{l\dagger} U^\nu \Phi$ mixing matrix is given now by

$$V_{MNS} \sim \begin{pmatrix} -0.9979 & -0.0461i & 0.0445i & 0.0095 \\ -0.0630 & 0.8544i & -0.5048i & -0.1056 \\ -0.0151 & -0.5175i & -0.8363i & -0.1807 \\ -0.0001 & 0.0029i & 0.2094i & -0.9778 \end{pmatrix}. \quad (20)$$

Notice that the 3×3 (non-hermitian) submatrix involving the known charged leptons and the three active neutrinos is again of the SMA solution type of the solar neutrino problem [21,23]. However see the discussion above. In this case we also obtain the mass square differences which are appropriate for the solar and atmospheric neutrino data. From the fourth raw and column we see that the heavy charged lepton couples mainly with the fourth neutrino and the usual charged leptons couple weakly to this neutrino. Notwithstanding, also in this case the mixing matrix is not orthogonal and it is not possible to make a direct comparison with the usual 3×3 fitting of the data and, as we said before, using a different set of parameters in Eqs. (5) and (18) we should obtain different mixing matrices.

Then, in this context we have four phenomenological neutrinos denoted by $\nu_{eL}, \nu_{\mu L}, \nu_{\tau L}, (N'^c)_L$ in the symmetry basis. So, in general, we have the possibility of implementing “2+2” [28], “3+1” [29–31], or intermediate neutrino mixing schemes. However it is worth to recall that it is not clear yet that we need four neutrinos since other type of solutions like flavor changing neutrino interactions are still possible for the solar neutrino [32,33] and also for the case of atmospheric neutrino, more exotic solutions like neutrino decay and decoherence are not ruled out yet [34], and moreover, several of these mechanisms can be working at the same time. Notice also that without the sextet we have no chance to assume that the charged lepton mass matrix is diagonal, and that the fourth row gives the couplings of the neutrinos with the heavy charged lepton which are very small.

Finally, we would like to remember that in the present model there are energy independent flavor conversions [35].

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